

**$\lambda$ -prophage induction modeled as a cooperative failure mode of lytic repression**Nicholas Chia,<sup>1,2</sup> Ido Golding,<sup>2</sup> and Nigel Goldenfeld<sup>1,2</sup><sup>1</sup>*Institute for Genomic Biology, University of Illinois at Urbana–Champaign, 1206 West Gregory Drive, Urbana, Illinois 61801, USA*<sup>2</sup>*Loomis Laboratory of Physics and Center for the Physics of Living Cells, University of Illinois at Urbana–Champaign,**1110 West Green Street, Urbana, Illinois 61801, USA*

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We analyze a system-level model for lytic repression of  $\lambda$  phage in *E. coli* using reliability theory, showing that the repressor circuit comprises four redundant components whose failure mode is prophage induction. Our model reflects the specific biochemical mechanisms involved in regulation, including long-range cooperative binding, and its detailed predictions for prophage induction in *E. coli* under ultraviolet radiation are in good agreement with experimental data.

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**I. INTRODUCTION**

When some phages infect a bacterial cell, there can be two possible outcomes or pathways [1]. In the lytic pathway, the phage hijacks the cell's machinery to replicate itself many times and to release the replicates by breaking open or lysing the cell. In the lysogenic pathway, the phage integrates its genome into that of the host microbe, becoming a prophage, but otherwise does not damage the cell. This lysogenic state is very stable [2]; however, an insult to the cell through, for example, starvation or exposure to ultraviolet (UV) radiation [3] can trigger a process known as prophage induction [4]: the prophage is excised from the cell's genome, and viral replication occurs leading to cell lysis. The most well-studied lysogenic system is the bacteriophage  $\lambda$ , or  $\lambda$  phage, which infects *Escherichia coli*. Understanding the lysis-lysogeny system in detail is important, because this system is one of the simplest examples of a gene regulatory network [5]—a pervasive and fundamental form of biological organization and function, whose principles are still being elucidated. Although there has been considerable interest recently in the role of stochasticity [6] in the switching behavior between lytic and lysogenic states as part of the phage lifecycle [7–11], here, we focus on UV prophage induction, where a different mechanism is involved.

UV prophage induction experiments exhibit threshold behavior [12], in which the fraction of induced lysogens, (i.e., prophage containing cells) rapidly increases as a function of the UV dose. Under typical laboratory growth conditions, the fraction of induced lysogens versus the UV dosage obeys a power law with a power very close to 4 [12]. Power-law behaviors of this type can arise in several ways: (i) as an event caused by four independent hits on a “target” (target theory [13]) or (ii) a chemical equilibrium reaction involving a substrate bound to four chemical species and quantified by the empirical Hill equation for chemical kinetics [14,15]. However, with little connection to the biochemical regulatory mechanisms of  $\lambda$ -phage lytic repression [4,16], such approaches are unable to make predictions *a priori* for power-law behavior and values of exponents. Another alternative perspective is to view UV induction in the framework of the standard stochastic model of lytic repression [8,9] with adjusted rate constants.

In this Rapid Communication, we connect the role of DNA loops and long-range cooperative binding [17,18] in the biochemical picture of lytic repression to the phenomenology of prophage induction. Our approach is to abstract the biochemistry into a system-level description, in which the lytic repressor circuit is represented as a device comprised of a number of redundant elements and one failure mode, lysis. This allows us to draw connections between the biochemical regulatory mechanism and the reliability theory [19–21] and also predict the characteristic power law for UV prophage induction.

**II. BIOCHEMISTRY OF LYTIC REPRESSION**

Figure 1 illustrates a widely accepted model of the lytic repressor switch in  $\lambda$  phage (for a review, see [1]). The lytic repressor molecule CI dimerizes and binds to specific DNA sites in the  $O_L$  and  $O_R$  control boxes,  $O_{L1}$  and  $O_{L2}$ , blocking

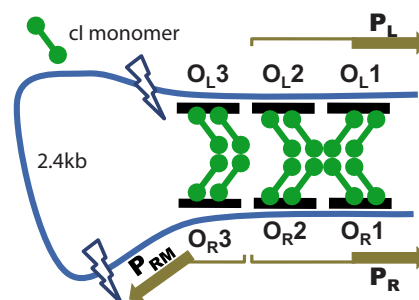


FIG. 1. (Color online) Schematic of the  $\lambda$ -phage lytic repression system. Attachment of RNA polymerase to promoter regions, indicated here by arrows, lead to gene expression.  $P_R$  and  $P_L$  lead to expression of genes in the lytic pathway. However, CI dimer binding at  $O_{R1}$  or  $O_{R2}$  blocks transcription of  $P_R$ , while CI binding at  $O_{L1}$  and  $O_{L2}$  blocks transcription of  $P_L$ .  $O_{R3}$  likewise regulates the transcription of genes by the promoter region  $P_{RM}$  while CI bound to  $O_{R2}$  promotes transcription of genes from  $P_{RM}$ . Dimers of CI are capable of forming stable quadramers when attached to adjacent sites such as at  $O_{R1}$  and  $O_{R2}$  [7,8] (modeled by [22]). Furthermore,  $O_{R1}$ - $O_{R2}$  and  $O_{L1}$ - $O_{L2}$  quadramers can form a stable octamer in a long-range interaction typically spanning 2.4 kb, but up to 3.8 kb [23,24].

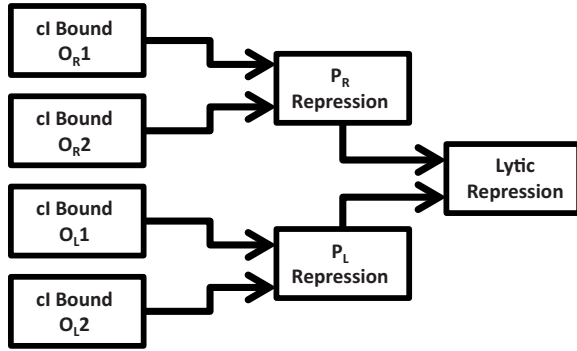


FIG. 2.  $\lambda$ -phage lytic repression regulatory circuit. CI bound to  $O_{R1}$  or  $O_{R2}$  blocks expression of lytic genes under the control of  $P_R$ . CI bound to  $O_{L1}$  or  $O_{L2}$  blocks expression of lytic genes under the control of  $P_L$ . Since both sets of genes under the control of  $P_R$  and  $P_L$  are required, blocking expression of either effectively represses lysis. Thus, each of the four CI bound components blocks lysis when intact. These can be seen as redundant elements that perform the same task.

expression of genes under the control of the  $P_L$  promoter. In the  $O_R$  control box,  $O_{R1}$  and  $O_{R2}$  regulate  $P_R$ . Since only free  $O_{R1}$  and  $O_{R2}$  sites allow the expression of  $P_R$  controlled genes, CI binding at either  $O_{R1}$  or  $O_{R2}$  suffices for repression of  $P_R$ . Figure 2 sketches the relationship between these two  $O_R$  sites and  $P_R$  expression. The same applies to the role of  $O_{L1}$  and  $O_{L2}$  in suppressing  $P_L$ . Also, as drawn in Fig. 2, the expression of genes under the control of both  $P_L$  and  $P_R$  promoters lead to lytic development of the prophage [25]. Derepression of all of the four binding sites results in lysis while bound CI dimers at any site block the lytic pathway. In UV induction, RecA-mediated autocleavage of CI monomers deprives the binding sites of available CI dimers. RecA-mediated autocleavage can happen once RecA is activated as part of the host SOS response to DNA damage [26].

### III. ABSTRACTION OF THE LYTIC REPRESSION CIRCUIT

Prophage induction can be understood as the failure of the lytic repression circuit, which consists of four redundant components that each prevents lysis. Each has a failure rate  $\mu_i$  ( $i=1, \dots, 4$ ) per UV dose  $x$  and a corresponding survival probability  $p_i = \exp(-\mu_i x)$ . Figure 2 diagrams the relevant aspects of the  $\lambda$ -phage lytic repressor regulatory system. Each component consists of a lytic repressor CI dimer bound to one of the four specific DNA sites, i.e.,  $O_{R1}$ ,  $O_{R2}$ ,  $O_{L1}$ , or  $O_{L2}$ . Each site regulates the expression of genes essential to the lytic pathway by its influence on either the promoter  $P_R$ , by  $O_{R1}$  and  $O_{R2}$ , or  $P_L$ , by  $O_{L1}$  and  $O_{L2}$ . Thus, these four components have redundant functionality, i.e., repressing lysis. Since the suppression of genes under the control of either promoter keeps lysis in check, only damage to the final component results in lysis.

### IV. RELIABILITY THEORY OF THE REPRESSOR CIRCUIT

To understand the failure rate of the system, i.e., the fraction of cells lysed, note that the probability of failure, for UV

dose  $x$ , is  $1 - p_i$  for each of the four redundant components (the four CI dimers bound to  $O_{R1}$ ,  $O_{R2}$ ,  $O_{L1}$ , and  $O_{L2}$ ) in the lytic repression system (see Fig. 2). In general, the probability of failure as a function of total UV dose,  $F$ , for a system of  $n$  redundant components is

$$F = \sum_{\Psi \in A} w(\Psi) \prod_{i=1}^n [1 - p_i(\Psi)], \quad (1)$$

where  $A$  represents the set of all combinations of possible paths leading to failure of all four components and  $w(\Psi)$  gives the probability of a particular combination of failures  $\Psi$  occurring. The probability of failure  $p_i(\Psi)$  accounts for the dependence of the failure of the  $i$ th component on the state of the repressor system, reflecting the role of cooperative bindings between CI dimers shown in Fig. 1. We model the effects of radiation on the failure rate of the lytic repression system, by assuming that near the threshold  $x_c$ ,  $p_i(\Psi) = \exp[-\mu_i(\Psi)(x - x_c) + O((x - x_c)^2)]$ .

By taking measurements of the fraction of failed systems and Eq. (1), the number of redundant elements in the system can be deduced. The fraction of failed units is then given by

$$F(x) = \sum_{\Psi \in A} w(\Psi) \prod_{i=1}^n \{1 - \exp[-\mu_i(\Psi)x]\} \approx (\mu x)^n \quad (2)$$

in the limit  $\mu x \ll 1$  and where  $\mu = \sum_{\Psi \in A} w(\Psi) [\prod_{i=1}^n \mu_i(\Psi)]^{1/n}$  is an effective failure rate. The UV prophage induction curve describes the fraction of cells lysed as a function of UV dose  $x$  and is predicted to follow Eq. (2) with  $n=4$ . In other words, the fraction of cells lysed can be computed from the effective failure rate  $\mu$  for the four CI dimer bindings at  $O_{R1}$ ,  $O_{R2}$ ,  $O_{L1}$ , and  $O_{L2}$ . Note that the power-law behavior arises from the cooperative CI-DNA interactions at these four sites while the effective rate of failure  $\mu$  varies between different experimental systems and depends on a number of parameters such as CI-CI dimer short- and long-range cooperative bindings [17,18] or varying levels of RecA activity [27,26]. Alternatively, mutant  $cI$  alleles may offer operator site binding affinities [28]. These changes in  $\mu$  shift the threshold of induction but do not affect the functional form of the UV dose curve itself. Below, we test Eq. (2) against data from experiments on radiative induced lysis and extract  $\mu$ .

### V. EXPERIMENTAL TESTS

In order to measure the fraction of lysogens induced as a function of UV dose, we followed standard protocols [29]. Briefly, exponentially growing lysogenic cells were harvested and resuspended in buffer. The cells were then irradiated by a germicidal UV lamp in dim ambient light for a range of doses at  $\sim 1$  (J/m<sup>2</sup>)/s. After irradiation, aliquots were diluted into growth medium, shaken for 2 h at 37 °C in the dark, treated with CHCl<sub>3</sub> and titered for plaque forming units.

In our abstraction of the lytic repression system, we noted four redundant components, as shown in Fig. 2. Since this implies  $n=4$ , we can test the applicability of Eq. (2) to our experimental results by plotting  $F/(x^4)$  against the UV dose  $x$ , as shown in Fig. 3.

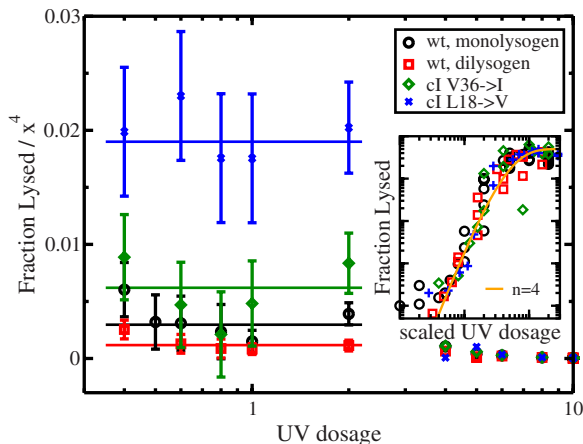


FIG. 3. (Color online) The UV prophage induction curve.  $F/(x^4)$  versus UV dosage  $x$ . Error bars were calculated from the sum of the residuals across UV dosages from 0.3 to 2 J/m<sup>2</sup>. The data are taken from different  $\lambda$  strains. Wild-type prophages were integrated as monolysogens and dilysogens with one or two inserted  $\lambda$  genomes, respectively. All four induction curves scale approximately as a power of 4 for small UV doses (0.2–2 J/m<sup>2</sup>). (Inset) Log-log plot of the same data. Curves are right and left shifted, so that they overlay each other. The line represents Eq. (1) for  $n=4$ .

In Fig. 3, the effective rate of failure  $\mu$  manifests itself as a horizontal line. As shown in Fig. 3, experimental data are consistent with Eq. (2) for a certain range of the UV dosage. Disagreement between theory and experiment occurs at both low and high UV dosages. The breakdown in both these regimes is readily interpreted. For extremely low dosage (near zero), spontaneous lysis events not induced by UV irradiation become the strongest contributing factor to the failure rate. These failures come from other events, such as spontaneous RecA activity [4] and mutations to  $\lambda$  phage [2]. At high radiation doses (not plotted, see [12]), the fraction of lysis does not saturate at 1 and instead begins falling with respect to the UV dose. Here, damage to the lytic pathway likely results in the inability of a cell to lyse, either because key components of the lytic pathway or host metabolism have been crucially damaged. In other words, this drop is not the effect of a lower failure rate of lytic repression, but a reflection of the high failure rate of other cellular systems upon which lysis relies. In these cases, failure of lytic repression cannot be detected by cell lysis since the lytic response has been disabled.

## VI. DISCUSSION

Our model postulates that induction, or a failure event, is dominated primarily by CI dimer dissociation. The mechanisms that relate UV dosage to operator site bound CI-dimer dissociation is unclear and an area of ongoing research [30,31]. In our model, UV radiation leads to DNA lesions which then result in the activation of RecA which then promotes the autocleavage of CI monomers from the cytoplasm. CI monomer depletion then leads to the depletion of available CI dimers for maintaining the repression of  $P_R$  and  $P_L$  as the rate of CI association falls below the constant disso-

TABLE I. List of predicted and observed changes in inducibility tabulated according to mutation, either in  $\lambda$  or in the *E. coli* host.  $\uparrow$  indicates that the change results an increase in  $\mu$ , while  $\downarrow$  indicates a decrease. Predictions based on reliability theory match with currently available data.

Strain	Phenotype	Theory	Data	Ref.
$\lambda$ cI <i>ind</i> <sup>s</sup> -1	Faster CI cleavage	$\uparrow$	$\uparrow$	[27]
<i>lexA51 recA441</i>	Increased RecA activity	$\uparrow$	$\uparrow$	[27]
$\lambda$ cI <i>ind</i> 543	Stronger CI dimerization	$\downarrow$	$\downarrow$	[27]
$\lambda$ cI Y210 $\rightarrow$ N	Disrupts CI dimer-dimer cooperative binding	$\uparrow$	$\uparrow$	[36]
$O_R2^*$	Weaker $O_R2$ -CI binding	$\uparrow$	$\uparrow$	[15] <sup>a</sup>

<sup>a</sup>Measured  $P_R$  expression.

ciation rate. Near the threshold of prophage induction, only the most rapidly varying parameter is important. In our model, it is the rapidly varying CI-dimer operator site bindings about the threshold point for lytic induction  $x_c$  that justifies the approximation made in Eq. (2). Accordingly we do not need to make any assumptions about the functional relationships between the many mechanisms in the path from UV radiation to CI monomer depletion, suggesting insensitivity of the lytic-lysogenic decision to these other mechanisms in the vicinity of the threshold.

Chemical kinetic models assume that since the time scale of CI dimer dissociation falls an order of magnitude below the time required for lysis, the CI bindings can be regarded as being adiabatically slaved to CI monomer concentration [8,9], and thus not a determining factor in the switch. The prophage induction curve and its power-law behavior then arise from the behavior of the CI monomer depletion. However, the measured dissociation rates, which range from  $\sim 30$  s [28] to  $\sim 4$  min [32], do not provide the separation of time scales necessary to justify the instantaneous dissociation rate approximation. Also, as can be seen from the data presented in Fig. 3, the power-law dependence  $n=4$  is robust and not sensitive to the different strains or experimental conditions. In our model, the power  $n$  reflects the number of redundant elements—not the properties of the individual components—and so is robust. This can be directly tested by manipulating the number of redundant operator sites and measuring the resultant power-law dependence. Reference [15] measured  $P_R$  expression as a function of CI concentration. As shown in Fig. 1,  $P_R$  expression is regulated by two bound operator sites, so we predict that  $P_R$  expression should be described by a kinetic curve with an  $n=2$  power-law dependence, as has been previously noted [8]. The kinetic data of Ref. [15] are indeed consistent with the prediction we have made here, based on our system-level abstraction of the underlying biology. This finding supports our view that the threshold behavior of prophage induction is determined by CI, and not by other steps in the repressor circuit. To establish this more conclusively, it would be necessary to check that  $P_L$  expression also exhibits the predicted  $n=2$  power-law behavior.

In our model, the failure events are dominated by CI dissociation with Cro providing the bistability by not allowing

newly produced CI dimers to re-engage the operator binding sites that would repress  $P_R$  and  $P_L$ . This implies that the repressor sites are the key to stabilizing the lysogenic state while the presence of Cro does not play a role in the switching [33,34] except to enforce commitment to the developmental transition [35]. Our model also suggests a mechanism for abortive induction events that are sometimes observed [4]. In our model these arise when unblocked  $P_L$  transcribes the genes required for excision (see [1]) while  $P_R$  remains blocked.

As shown by Fig. 3, different conditions lead to different values for  $\mu$ . We can use reliability theory to anticipate the trends in variation of  $\mu$  between two similar experiments. The rate of component failure depends on a number of variables including RecA activity, CI concentration, binding strength of repressor sites, and stability of CI to autocleavage. Here, damage to redundant components corresponds to

the dissociation of CI dimers from  $O_L1/2$  and  $O_R1/2$  sites (see Fig. 2). Table I lists theoretical predictions for the variation in the failure rate of the lytic repressor or inducibility arising from possible laboratory manipulations of the rate of failure  $\mu$ .

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- [1] M. Ptashne, *A Genetic Switch: Phage Lambda Revisited* (Cold Spring Harbor Laboratory Press, New York, 2004).
- [2] K. Bæk, S. Svenningsen, H. Eisen, K. Sneppen, and S. Brown, *J. Mol. Biol.* **334**, 363 (2003).
- [3] S. Luria and R. Latarjet, *J. Bacteriol.* **53**, 149 (1947).
- [4] A. Oppenheim, O. Kobiler, J. Stavans, and D. Court, *Annu. Rev. Genet.* **39**, 409 (2005).
- [5] J. Hasty *et al.*, *Nat. Rev. Genet.* **2**, 268 (2001).
- [6] A. Arkin, J. Ross, and H. McAdams, *Genetics* **149**, 1633 (1998).
- [7] G. Ackers, A. Johnson, and M. Shea, *Proc. Natl. Acad. Sci. U.S.A.* **79**, 1129 (1982).
- [8] M. Shea and G. Ackers, *J. Mol. Biol.* **181**, 211 (1985).
- [9] E. Aurell, S. Brown, J. Johanson, and K. Sneppen, *Phys. Rev. E* **65**, 051914 (2002).
- [10] A. Lipshtat, A. Loinger, N. Q. Balaban, and O. Biham, *Phys. Rev. Lett.* **96**, 188101 (2006).
- [11] W. Bialek and S. Setayeshgar, *Phys. Rev. Lett.* **100**, 258101 (2008).
- [12] H. Kneser, *Virology* **28**, 701 (1966).
- [13] G. Stent and J. Dohm, *Molecular Biology of Bacterial Viruses* (WH Freeman, San Francisco, 1963).
- [14] A. Hill, *J. Physiol. (London)* **40**, 4 (1910).
- [15] N. Rosenfeld, J. Young, U. Alon, P. Swain, and M. Elowitz, *Science* **307**, 1962 (2005).
- [16] D. Court, A. Oppenheim, and S. Adhya, *J. Bacteriol.* **189**, 298 (2007).
- [17] I. Dodd, K. Shearwin, A. Perkins, T. Burr, A. Hochschild, and J. Egan, *Genes Dev.* **18**, 344 (2004).
- [18] I. Dodd, K. Shearwin, and K. Sneppen, *J. Mol. Biol.* **369**, 1200 (2007).
- [19] I. Gertsbakh, *Adv. Appl. Probab.* **16**, 147 (1984).
- [20] L. Gavrilov and N. Gavrilova, *J. Theor. Biol.* **213**, 527 (2001).
- [21] L. Gavrilov and N. Gavrilova, in *Handbook of Models for Human Aging*, edited by P. Conn (Elsevier Academic Press, Burlington, MA, 2006), pp. 45–68.
- [22] J. M. G. Vilar and L. Saiz, *Phys. Rev. Lett.* **96**, 238103 (2006).
- [23] B. Revet, B. von Wilcken-Bergmann, H. Bessert, A. Barker, and B. Muller-Hill, *Curr. Biol.* **9**, 151 (1999).
- [24] I. Dodd, A. Perkins, D. Tsemitsidis, and J. Egan, *Genes Dev.* **15**, 3013 (2001).
- [25] I. Dodd, K. Shearwin, and J. Egan, *Curr. Opin. Genet. Dev.* **15**, 145 (2005).
- [26] J. Little and D. Mount, *Cell* **29**, 11 (1982).
- [27] M. Dutreix, A. Bailone, and R. Devoret, *J. Bacteriol.* **161**, 1080 (1985).
- [28] H. Nelson and R. Sauer, *Cell* **42**, 549 (1985).
- [29] J. Little, D. Shepley, and D. Wert, *EMBO J.* **18**, 4299 (1999).
- [30] J. Little, *Proc. Natl. Acad. Sci. U.S.A.* **102**, 5310 (2005).
- [31] V. Galkin, X. Yu, J. Bielnicki, D. Ndjonka, C. Bell, and E. Egelman, *J. Mol. Biol.* **385**, 779 (2008).
- [32] Y. Wang, L. Guo, I. Golding, E. Cox, and N. Ong, *Biophys. J.* **96**, 609 (2009).
- [33] S. Svenningsen, N. Costantino, D. Court, and S. Adhya, *Proc. Natl. Acad. Sci. U.S.A.* **102**, 4465 (2005).
- [34] S. Atsumi and J. Little, *Proc. Natl. Acad. Sci. U.S.A.* **103**, 4558 (2006).
- [35] R. Schubert, I. Dodd, J. Egan, and K. Shearwin, *Genes Dev.* **21**, 2461 (2007).
- [36] A. Babic and J. Little, *Proc. Natl. Acad. Sci. U.S.A.* **104**, 17741 (2007).